

Finite length Solenoid potential and field

The origin document was written in 2011. There are sign error and factor 2 error at that time. I update and redo the calculation in 2021. It is a supplementary document for the webpage <https://nukephysik101.wordpress.com/2011/07/17/the-magnetic-field-of-a-finite-length-solenoid/>. People should also refer to the calculation of the single coil. At the time of writing this document, I was a PhD student at Tokyo University.

Tsz Leung Tang, Saturday, July 10, 2021

The surface current density is (Jackson, 1998):

$$\vec{K} = \frac{I}{L} \delta(\rho - a) (-\sin \phi, \cos \phi, 0), \quad z \in \left(-\frac{L}{2}, \frac{L}{2}\right).$$

In above, L is the total length of the solenoid, I is the current, (ρ, ϕ, z) are the cylindrical coordinate, and a is the radius of the solenoid. The general form of vector potential is

$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3x'.$$

By symmetry, the vector potential is azimuthal.

$$\vec{A} = A_\phi \hat{\phi} = \hat{\phi} \frac{\mu_0 I}{4\pi L} \int_{-\frac{L}{2}}^{\frac{L}{2}} \int_0^{2\pi} \int_0^\infty \frac{\delta(\rho' - a) \cos \phi'}{\sqrt{\rho^2 + \rho'^2 + (z - z')^2 - 2a\rho \cos \phi'}} \rho' d\rho' d\phi' dz',$$
$$A_\phi = \frac{\mu_0 I a}{2\pi L} \int_{-\frac{L}{2}}^{\frac{L}{2}} \int_0^\pi \frac{\cos \phi'}{\sqrt{\rho^2 + a^2 + (z - z')^2 - 2a\rho \cos \phi'}} d\phi' dz'.$$

Simplify the form by setting $\zeta = (z - z')$ and the integration of ζ is a log function (Edmund E. Callaghan, 1960), we have,

$$\int_0^\pi \cos \phi' \left[\ln \left(\zeta + \sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'} \right) \right]_{\zeta_-}^{\zeta_+} d\phi' = \left[\int_0^\pi \cos \phi' \ln(\zeta + \alpha(\zeta)) d\phi' \right]_{\zeta_-}^{\zeta_+},$$

$$\alpha(\zeta) = \sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'}, \quad \zeta_\pm = z \mp \frac{L}{2}.$$

Then,

$$A_\phi = \frac{\mu_0 I a}{2\pi L} \left[\int_0^\pi \cos \phi' \ln(\zeta + \alpha(\zeta)) d\phi' \right]_{\zeta_-}^{\zeta_+}$$

Integration by path gives,

$$\int_0^\pi \cos \phi' \ln(\zeta + \alpha(\zeta)) d\phi' = \sin \phi' \ln(\zeta + \alpha(\zeta)) \Big|_0^{2\pi} - \int_0^\pi \sin \phi' d(\ln(\zeta + \alpha(\zeta)))$$

The first term is zero because of the $\sin \phi'$, and the derivative of $\ln(\zeta + \alpha(\zeta))$ is:

$$\frac{d \ln(\zeta + \alpha(\zeta))}{d\phi'} = \frac{\rho a \sin \phi'}{(\alpha(\zeta) + \zeta)\alpha(\zeta)}$$

Multiple by $(\alpha(\zeta) - \zeta)/(\alpha(\zeta) - \zeta)$

$$\begin{aligned} &= \frac{(\alpha(\zeta) - \zeta)\rho a \sin \phi'}{(\alpha^2(\zeta) - \zeta^2)\alpha(\zeta)} \\ &= \frac{\rho a \sin \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi')} - \frac{\zeta \rho a \sin \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi')\alpha(\zeta)} \end{aligned}$$

The first term is a constant of ζ , and $[x]_{\zeta_-}^{\zeta_+} = 0$, then,

$$\begin{aligned} \int_0^\pi \sin \phi' d(\ln(\zeta + \alpha(\zeta))) &= - \int_0^\pi \frac{\zeta \rho a \sin^2 \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi')\alpha(\zeta)} d\phi' \\ A_\phi &= \frac{\mu_0 I a^2 \rho}{2\pi L} \left[\zeta \int_0^\pi \frac{\sin^2 \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi')\sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'}} d\phi' \right]_{\zeta_-}^{\zeta_+} \end{aligned}$$

Now, since the $\cos \phi$ is the same as interval $(-\pi, 0)$ and $(0, \pi)$, we can change the sign, and replacing $\phi = 2\theta$, using $\cos(2\theta) = 1 - 2 \sin \theta$, we have

$$\begin{aligned} &\int_0^\pi \frac{\sin^2 \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi')\sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'}} d\phi' \\ &= \int_0^\pi \frac{\sin^2 \phi'}{(\rho^2 + a^2 + 2\rho a \cos \phi')\sqrt{\zeta^2 + \rho^2 + a^2 + 2\rho a \cos \phi'}} d\phi' \\ &= 2 \int_0^{\frac{\pi}{2}} \frac{\sin^2(2\theta)}{((a + \rho)^2 - 4\rho a \sin^2 \theta)\sqrt{\zeta^2 + (a + \rho)^2 - 4\rho a \sin^2 \theta}} d\theta \\ &= \frac{kh^2}{4(\sqrt{a\rho})^3} \int_0^{\frac{\pi}{2}} \frac{\sin^2(2\theta)}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\ &= \frac{kh^2}{(\sqrt{a\rho})^3} \int_0^{\frac{\pi}{2}} \frac{\sin^2 \theta - \sin^4 \theta}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta, \quad \sin^2(2\theta) = 4 \sin^2 \theta - 4 \sin^4 \theta \\ h^2 &= \frac{4a\rho}{(a + \rho)^2} \\ k^2 &= \frac{4a\rho}{(a + \rho)^2 + \zeta^2} \end{aligned}$$

In the above calculation, in the 1st step, the $\cos \phi$ denominator change sign means the while function flipped horizontally on the $\phi = \frac{\pi}{2}$.

The integral can be spitted into 2 parts, the first part is (Milton Abramowitz, 1965) (NIST Digital Library of Mathematical Functions) :

$$\begin{aligned}
 & \int_0^{\frac{\pi}{2}} \frac{\sin^2 \theta}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= -\frac{1}{h^2} \int_0^{\frac{\pi}{2}} \frac{1 - h^2 \sin^2 \theta - 1}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= -\frac{1}{h^2} \int_0^{\frac{\pi}{2}} \frac{1}{\sqrt{1 - k^2 \sin^2 \theta}} - \frac{1}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= \frac{1}{h^2} (\Pi(h^2, k^2) - K(k^2))
 \end{aligned}$$

Here

$$\Pi(n, m) = \int_0^{\frac{\pi}{2}} \frac{1}{(1 - n \sin^2 \theta)\sqrt{1 - m \sin^2 \theta}} d\theta$$

It is the elliptic integral of 3rd kind.

The 2nd part is:

$$\begin{aligned}
 & \int_0^{\frac{\pi}{2}} \frac{-\sin^4 \theta}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= \frac{1}{h^4} \int_0^{\frac{\pi}{2}} \frac{1 - h^4 \sin^4 \theta - 1}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= \frac{1}{h^4} \int_0^{\frac{\pi}{2}} \frac{1 + h^2 \sin^2 \theta}{\sqrt{1 - k^2 \sin^2 \theta}} - \frac{1}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
 &= \frac{1}{h^4} \left(K(k^2) + \frac{h^2}{k^2} (K(k^2) - E(k^2)) - \Pi(h^2, k^2) \right)
 \end{aligned}$$

Thus, combine everything and we have,

$$A_\phi = \frac{\mu_0 I}{2\pi L} \frac{1}{\sqrt{\rho}} \left[\zeta k \left(\frac{k^2 + h^2 - h^2 k^2}{h^2 k^2} K(k^2) - \frac{1}{k^2} E(k^2) + \frac{h^2 - 1}{h^2} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

The magnetic field is the curl

$$\begin{aligned}
 B_\rho &= [\nabla \times A_\phi]_\rho = -\frac{\partial}{\partial z} (A_\phi) = -\frac{\partial}{\partial \zeta} (A_\phi) \\
 B_z &= [\nabla \times A_\phi]_z = \frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho A_\phi) = \frac{1}{\rho} A_\phi + \frac{\partial A_\phi}{\partial \rho}
 \end{aligned}$$

Using the derivative formulae for elliptic integral:

$$\begin{aligned}\frac{d}{dk}K(k^2) &= -\frac{1}{k}K(k^2) + \frac{1}{k(1-k^2)}E(k^2) \\ \frac{d}{dk}E(k^2) &= -\frac{1}{k}K(k^2) + \frac{1}{k}E(k^2) \\ \frac{d}{dk}\Pi(h^2, k^2) &= \frac{-k}{(1-k^2)(h^2-k^2)}E(k^2) - \frac{k}{(h^2-k^2)}\Pi(v^2, x^2)\end{aligned}$$

Since derivative of ζ and ρ is through derivative of k , we compute,

$$\begin{aligned}\frac{d}{dk}\left(k\left(\frac{k^2+h^2-h^2k^2}{h^2k^2}K(k^2) - \frac{1}{k^2}E(k^2) + \frac{h^2-1}{h^2}\Pi(h^2, k^2)\right)\right) \\ = -\frac{1}{k^2}K(k^2) + \frac{h^2}{k^2(h^2-k^2)}E(k^2) + \frac{h^2-1}{(h^2-k^2)}\Pi(h^2, k^2)\end{aligned}$$

And,

$$\begin{aligned}\frac{dk}{dz} &= -\frac{k^3\left(z \pm \frac{L}{2}\right)}{4a\rho}, \quad \frac{dk}{d\rho} = \frac{k}{2\rho} - \frac{k^3(a+\rho)}{4a\rho} \\ \frac{d}{dz}(zf(z)) &= f(z) + z\frac{df(z)}{dz}, \quad \frac{d}{d\rho}\left(\frac{1}{\sqrt{\rho}}f(\rho)\right) = -\frac{f(\rho)}{2\sqrt{\rho^3}} + \frac{1}{\sqrt{\rho}}\frac{df(\rho)}{d\rho}\end{aligned}$$

Then

$$B_\rho = \frac{\mu_0 I}{2\pi L} \sqrt{\frac{a}{\rho}} \left[\left(\frac{k^2-2}{k}K(k^2) + \frac{2}{k}E(k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

Or, by using integration identity

$$\frac{d}{dx} \int_a^b f(x) dx = f(b) - f(a)$$

And the original formula of A_ϕ

$$A_\phi = \frac{\mu_0 I a}{2\pi L} \int_{-\frac{L}{2}}^{\frac{L}{2}} \int_0^\pi \frac{\cos \phi'}{\sqrt{\rho^2 + a^2 + (z-z')^2 - 2a\rho \cos \phi'}} d\phi' dz'.$$

We have,

$$B_\rho = \frac{\mu_0 I a}{2\pi L} \left[\int_0^\pi \frac{\cos \phi'}{\sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'}} d\phi' \right]_{\zeta_-}^{\zeta_+}$$

Then, we can verify the two expression is the same. Our result for B_ρ is

$$B_\rho = \frac{\mu_0 I}{2\pi L} \sqrt{\frac{a}{\rho}} \left[\left(\frac{k^2 - 2}{k} K(k^2) + \frac{2}{k} E(k^2) \right) \right]_{\zeta_-}^{\zeta_+}, \quad k^2 = \frac{4a\rho}{(a + \rho)^2 + \zeta^2}$$

To compute the z component, we need to be careful the $\Pi(h^2, k^2)$, since h^2 also contains ρ . And the z-component is:

$$B_z = \frac{\mu_0 I}{2\pi L} \frac{1}{2\sqrt{a\rho}} \left[\zeta k \left(K(k^2) + \frac{a - \rho}{a + \rho} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

To verify the B_z , we can compute $\frac{\partial A_\phi}{\partial \rho}$ from

$$A_\phi = \frac{\mu_0 I a}{2\pi L} \left[\int_0^\pi \cos \phi' \ln(\zeta + \alpha(\zeta)) d\phi' \right]_{\zeta_-}^{\zeta_+}$$

By

$$\frac{\partial}{\partial \rho} \ln(\zeta + \alpha(\zeta)) = \frac{\rho - a \cos \phi'}{\alpha(\zeta)(\alpha(\zeta) + \zeta)}$$

Using the same trick

$$\begin{aligned} \frac{\rho - a \cos \phi'}{\alpha(\zeta)(\alpha(\zeta) + \zeta)} &= \frac{(\rho - a \cos \phi')(\alpha(\zeta) - \zeta)}{\alpha(\zeta)(\alpha^2(\zeta) + \zeta^2)} \\ &= \frac{(\rho - a \cos \phi')(\alpha(\zeta) - \zeta)}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} = \frac{(\rho - a \cos \phi')}{(\rho^2 + a^2 - 2\rho a \cos \phi')} - \frac{(\rho - a \cos \phi')\zeta}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} \end{aligned}$$

Therefore:

$$\frac{\partial A_\phi}{\partial \rho} = \frac{\mu_0 I a}{2\pi L} \left[\int_0^\pi \frac{-\zeta \rho \cos \phi' + a\zeta \cos^2 \phi'}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} d\phi' \right]_{\zeta_-}^{\zeta_+}$$

Combined with

$$\frac{1}{\rho} A_\phi = \frac{\mu_0 I a}{2\pi L} \left[\int_0^\pi \frac{a\zeta \sin^2 \phi'}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} d\phi' \right]_{\zeta_-}^{\zeta_+}$$

Then the magnetic field is

$$B_z = \frac{\mu_0 I a}{2\pi L} \left[\zeta \int_0^\pi \frac{a - \rho \cos \phi'}{\alpha(\zeta)(\rho^2 + a^2 - 2\rho a \cos \phi')} d\phi' \right]_{\zeta_-}^{\zeta_+}$$

With the same trick,

$$\int_0^\pi \frac{a - \rho \cos \phi'}{(\rho^2 + a^2 - 2\rho a \cos \phi') \sqrt{\zeta^2 + \rho^2 + a^2 - 2\rho a \cos \phi'}} d\phi'$$

$$\begin{aligned}
&= \int_0^\pi \frac{a + \rho \cos \phi'}{(\rho^2 + a^2 + 2\rho a \cos \phi')\sqrt{\zeta^2 + \rho^2 + a^2 + 2\rho a \cos \phi'}} d\phi' \\
&= 2 \int_0^{\frac{\pi}{2}} \frac{a + \rho \cos(2\theta)}{((a + \rho)^2 - 4\rho a \sin^2 \theta)\sqrt{\zeta^2 + (a + \rho)^2 - 4\rho a \sin^2 \theta}} d\theta \\
&= \frac{kh^2}{4(\sqrt{a\rho})^3} \int_0^{\frac{\pi}{2}} \frac{(a + \rho) - 2\rho \sin^2 \theta}{(1 - h^2 \sin^2 \theta)\sqrt{1 - k^2 \sin^2 \theta}} d\theta \\
&= \frac{kh^2}{4(\sqrt{a\rho})^3} \left((a + \rho) \Pi(h^2, k^2) - \frac{2\rho}{h^2} (\Pi(h^2, k^2) - K(k^2)) \right) \\
&= \frac{k}{4(\sqrt{a\rho})^3} ((h^2(a + \rho) - 2\rho)\Pi(h^2, k^2) + 2\rho K(k^2)) \\
&= \frac{k}{4(\sqrt{a\rho})^3} \left(2\rho K(k^2) + \frac{2\rho(a - \rho)}{a + \rho} \Pi(h^2, k^2) \right) \\
&= \frac{k}{2a\sqrt{a\rho}} \left(K(k^2) + \frac{(a - \rho)}{a + \rho} \Pi(h^2, k^2) \right)
\end{aligned}$$

Thus we get the same result.

$$B_z = \frac{\mu_0 I}{2\pi L} \frac{1}{2\sqrt{a\rho}} \left[\zeta k \left(K(k^2) + \frac{a - \rho}{a + \rho} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+}$$

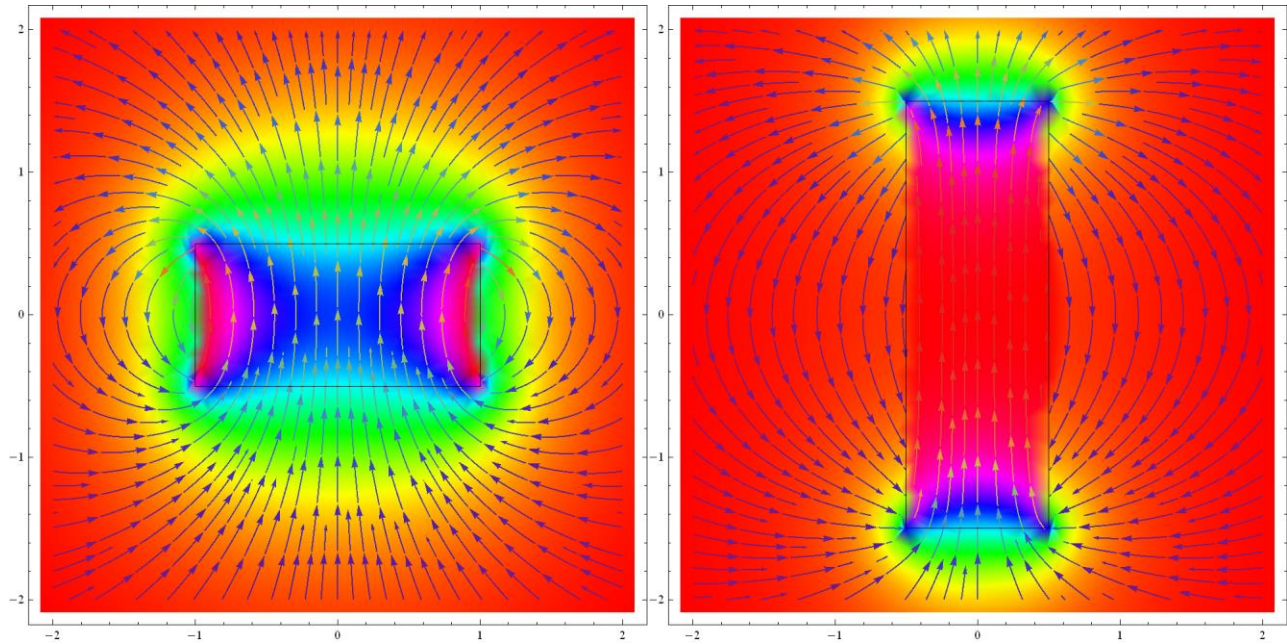
In conclusion, the field is defined by:

$$\begin{aligned}
A_\phi &= \frac{\mu_0 I}{2\pi L} \frac{1}{\sqrt{\rho}} \left[\zeta k \left(\frac{k^2 + h^2 - h^2 k^2}{h^2 k^2} K(k^2) - \frac{1}{k^2} E(k^2) + \frac{h^2 - 1}{h^2} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+} \\
B_\rho &= \frac{\mu_0 I}{2\pi L} \frac{1}{\sqrt{\rho}} \left[\left(\frac{k^2 - 2}{k} K(k^2) + \frac{2}{k} E(k^2) \right) \right]_{\zeta_-}^{\zeta_+} \\
B_z &= \frac{\mu_0 I}{2\pi L} \frac{1}{2\sqrt{a\rho}} \left[\zeta k \left(K(k^2) + \frac{a - \rho}{a + \rho} \Pi(h^2, k^2) \right) \right]_{\zeta_-}^{\zeta_+}
\end{aligned}$$

With

$$h^2 = \frac{4a\rho}{(a + \rho)^2}, \quad k^2 = \frac{4a\rho}{(a + \rho)^2 + \zeta^2}, \quad \zeta_\pm = z \pm \frac{L}{2}$$

Here are some plots:



Works Cited

- Edmund E. Callaghan, S. H. (1960). *The Magnetic Field of a Finite Solenoid (Technical note D-465)*. Washington, USA: Nation Aeronautics and Space Administration.
- Jackson, J. D. (1998). *Classical Electrodynamics*. John Wiley & Sons, Inc.
- Milton Abramowitz, I. A. (1965). *Handbook of mathematical functions : with formulas, graphs, and mathematical tables*. Dover.
- NIST Digital Library of Mathematical Functions*. (n.d.). Retrieved from <http://dlmf.nist.gov/>